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Modern cosmology 3: The Growth of Structure

- Growth of structure in an expanding universe
- The Jeans length
- Dark matter
- Large scale structure simulations
 effect of cosmological parameters
- Large scale structure data
 - ► galaxy surveys
 - ▶ cosmic microwave background

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Growth of structure

- Consider sphere of radius *R* and mass *M*
 - ▶ add small amount of mass so that $\rho = \overline{\rho}(1 + \delta(t))$
 - ▶ gravitational acceleration at edge of sphere is

$$\ddot{R} = -\frac{GM}{R^2} = -\frac{4\pi}{3}G\overline{\rho}R(1+\delta(t))$$

- ► since $\rho(t) = \rho_0 a^{-3}$, conservation of mass gives $R(t) \propto a(t)[1 + \delta(t)]^{-1/3}$
- ► differentiating this twice gives

$$\frac{\ddot{R}}{R} \approx \frac{\ddot{a}}{a} - \frac{\ddot{\delta}}{3} - \frac{2}{3}\frac{\dot{a}}{a}\dot{\delta}$$

Growth of structure

• Now have two equations for \ddot{R}/R

 $\frac{\ddot{R}}{R} = -\frac{GM}{R^3} = -\frac{4\pi}{3}G\overline{\rho}(1+\delta(t))$ $\frac{\ddot{R}}{R} \approx \frac{\ddot{a}}{a} - \frac{\ddot{\delta}}{3} - \frac{2}{3}\frac{\dot{a}}{a}\dot{\delta}$

Combine to get

$$-\frac{4}{3}\pi G\overline{\rho} - \frac{4}{3}\pi G\overline{\rho}\delta = \frac{\ddot{a}}{a} - \frac{1}{3}\ddot{\delta} - \frac{2}{3}\dot{\delta}\frac{\dot{a}}{a}$$

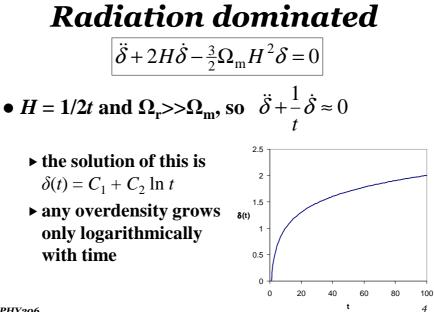
• Subtract off $\delta = 0$ case to get

$$\ddot{\delta} + 2H\dot{\delta} - \frac{3}{2}\Omega_{\rm m}H^2\delta = 0 \qquad \text{since} \qquad \Omega_{\rm m} = \frac{8\pi G\overline{\rho}}{3H^2}$$

R

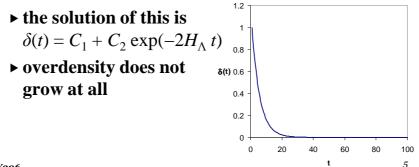
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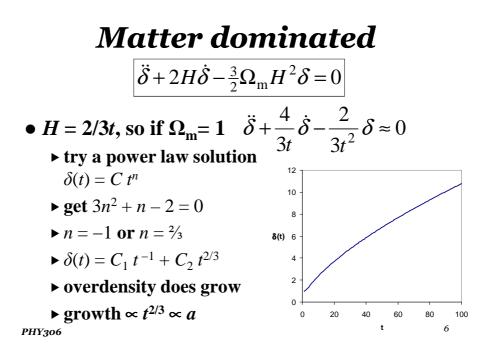
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Inflation/A dominated $\ddot{\delta} + 2H\dot{\delta} - \frac{3}{2}\Omega_{\rm m}H^2\delta = 0$

• $H = H_{\Lambda}$ and $\Omega_{\Lambda} >> \Omega_{\rm m}$, so $\ddot{\delta} + 2H_{\Lambda}\dot{\delta} \approx 0$





The Jeans length

- Astrophysical objects are stabilised against collapse by pressure forces
 - ► pressure forces travel at the speed of sound, $c_s = c \ (dP/d\varepsilon)^{1/2} = w^{1/2} c$
 - collapse occurs if the size of the object is smaller than the *Jeans length* _____

$$\ell_{\rm J} = c_{\rm s} \sqrt{\frac{\pi c^2}{G\bar{\varepsilon}}} = c^2 \sqrt{\frac{w\pi}{G\bar{\varepsilon}}}$$

where $\overline{\epsilon}$ is the mean energy density

 (exact numerical factor depends on details of derivation)

► can also be expressed as *Jeans mass PHY306*

The Jeans length

• We know that for a flat matter-dominated universe $H = (8\pi G\varepsilon/3c^2)^{1/2}$

• so
$$\ell_{\rm J} = c^2 \sqrt{\frac{w\pi}{G\overline{\varepsilon}}} = \sqrt{w} \times 2\pi \sqrt{\frac{2}{3}} \frac{c}{H}$$

- At last scattering $H^2 = \Omega_{m0} H_0^2 / a^3$, giving $c/H \sim 0.2$ Mpc for $\Omega_{m0} = 0.3$
 - ▶ so for radiation $l_{\rm J}$ ~ 0.59 Mpc
 - ► ~ horizon size, so nothing below horizon size collapses
 - For matter w = kT/µc² = 2.3 × 10⁻¹⁰, so ℓ_J ~ 15 pc
 equivalent to current size ~ 15 kpc, a small galaxy

Conclusions (radiation-dominated)

- Below Jeans length (~horizon size), nothing collapses
 - ► supported by pressure
- Above Jeans length
 - ▶ fluctuations in *matter* grow only logarithmically
 - ▶ fluctuations in *radiation* grow $\propto a^2$
 - derivation similar to p3 gives $\ddot{\delta} + 2H\dot{\delta} 4\Omega_r H^2 \delta = 0$
 - ► solve as in p6 to get $\delta \propto t$
 - ▶ but $l_J \propto a^2$, so as universe expands fluctuations will "enter the horizon" and stop growing

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Conclusions (matter-dominated)

- Before matter-radiation decoupling, photons and baryons form a single gas
 - ▶ this is stabilised by radiation pressure
 - ► density perturbations cannot grow
- After decoupling, baryon gas has galaxyscale Jeans length
 - ▶ galaxy-sized objects start to grow $\propto a$
 - but it turns out that this is too slow to produce structures we see
 - (note that our derivation is only valid for small δ)

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Dark Matter

- In fact we know $\Omega_{baryon} \sim 0.04$ from light element abundances, whereas $\Omega_m \sim 0.25$ from cluster X-ray data
 - most matter is not made of baryons (non-baryonic dark matter)
 - ▶ this does not couple to photons
 - ► dark matter structures can grow from time of matter-radiation equality (z ~ 3200)
 - ▶ this is much more promising

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Structure formation with dark matter

- Assume dark matter consists of a weakly interacting particle with mass *m*_{dm}
 - ▶ such a particle will be relativistic until the temperature drops to $T_{\rm dm} \sim m_{\rm dm}c^2/3k$
 - ▶ if this is less than $T_{\rm rm} \sim 10^4$ K, the dark matter is hot; if more, it is cold
 - ► for hot dark matter the minimum scale for collapse is given by $[ct_{dm} = ct_{rm}(T_{dm}/T_{rm})^2] \times (T_{dm}/2.74)$
 - ▶ for cold dark matter it is dwarf galaxy sized

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Conclusions (dark matter)

• Hot dark matter candidate: neutrino with mass ~2 eV/c²

- ► $T_{\rm dm} = 7750 \ {\rm K}$
- ▶ minimum length scale (for t_{rm} ~ 50000 yr)
 ~ 70 Mpc
- ► supercluster sized
- Large structures form first

- Cold dark matter candidate: WIMP with mass ~100 GeV/c²
 - ► $T_{\rm dm} = 4 \times 10^{14} \, {\rm K}$
 - minimum length scale set by Jeans length for matter
 - ► small galaxy sized
- Small structures form first

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